# EST Problem 1: Selected Solutions

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#### Abstract

Here I give a step-by-step solution to Q8 and Q10 of Prof Tew's EST worksheet 1.

## Problem 8

Let's begin with a simple two-electron case to understand how terms drop out, then generalize. For two electrons, the Slater determinant is:

$$|\Psi\rangle = \sqrt{2!}\hat{A}|\phi_1\phi_2\rangle = \frac{1}{\sqrt{2}}(|\phi_1\phi_2\rangle - |\phi_2\phi_1\rangle) \tag{1}$$

Let's evaluate  $\rho(x)$  for this two-electron case:

$$\rho(x) = \langle \Psi | \left[ \delta(x - x_1) + \delta(x - x_2) \right] | \Psi \rangle$$
<sup>(2)</sup>

$$= \frac{1}{2} \left( \left\langle \phi_1 \phi_2 \right| - \left\langle \phi_2 \phi_1 \right| \right) \left[ \delta(x - x_1) + \delta(x - x_2) \right] \left( \left| \phi_1 \phi_2 \right\rangle - \left| \phi_2 \phi_1 \right\rangle \right)$$
(3)

Let's focus on  $\delta(x - x_1)$  terms first:

$$\frac{1}{2} \left[ \left\langle \phi_{1}\phi_{2} \middle| \delta(x-x_{1}) \middle| \phi_{1}\phi_{2} \right\rangle 
- \left\langle \phi_{1}\phi_{2} \middle| \delta(x-x_{1}) \middle| \phi_{2}\phi_{1} \right\rangle 
- \left\langle \phi_{2}\phi_{1} \middle| \delta(x-x_{1}) \middle| \phi_{1}\phi_{2} \right\rangle 
+ \left\langle \phi_{2}\phi_{1} \middle| \delta(x-x_{1}) \middle| \phi_{2}\phi_{1} \right\rangle \right]$$
(4)

(5)

We can now evaluate each term:

• Term 1:  $\langle \phi_1 \phi_2 | \delta(x - x_1) | \phi_1 \phi_2 \rangle$ =  $|\phi_1(x)|^2 \langle \phi_2 | \phi_2 \rangle = |\phi_1(x)|^2$  • Term 2:  $\langle \phi_1 \phi_2 | \delta(x - x_1) | \phi_2 \phi_1 \rangle$ 

$$=\phi_1^*(x)\phi_2(x)\left\langle\phi_2|\phi_1\right\rangle=0$$
(6)

due to orbital orthogonality

• Term 3:  $\langle \phi_2 \phi_1 | \delta(x - x_1) | \phi_1 \phi_2 \rangle$ 

$$=\phi_2^*(x)\phi_1(x)\langle\phi_1|\phi_2\rangle=0$$
(7)

also due to orbital orthogonality

• Term 4:  $\langle \phi_2 \phi_1 | \delta(x - x_1) | \phi_2 \phi_1 \rangle$ 

$$= |\phi_2(x)|^2 \langle \phi_1 | \phi_1 \rangle = |\phi_2(x)|^2$$
(8)

Following the same process for  $\delta(x - x_2)$  terms:

$$\frac{1}{2} \left[ \left\langle \phi_{1}\phi_{2} \middle| \, \delta(x-x_{2}) \middle| \phi_{1}\phi_{2} \right\rangle 
- \left\langle \phi_{1}\phi_{2} \middle| \, \delta(x-x_{2}) \middle| \phi_{2}\phi_{1} \right\rangle 
- \left\langle \phi_{2}\phi_{1} \middle| \, \delta(x-x_{2}) \middle| \phi_{1}\phi_{2} \right\rangle 
+ \left\langle \phi_{2}\phi_{1} \middle| \, \delta(x-x_{2}) \middle| \phi_{2}\phi_{1} \right\rangle \right]$$
(9)

Again, cross terms drop out due to orthogonality, leaving us with:

$$\frac{1}{2}(|\phi_2(x)|^2 + |\phi_1(x)|^2) \tag{10}$$

Adding all surviving terms:

$$\rho(x) = \frac{1}{2}(|\phi_1(x)|^2 + |\phi_2(x)|^2) + \frac{1}{2}(|\phi_2(x)|^2 + |\phi_1(x)|^2)$$
(11)

$$= |\phi_1(x)|^2 + |\phi_2(x)|^2 \tag{12}$$

### Generalization to N Electrons

For the N-electron case, the same pattern holds but with more terms. The key points remain:

- 1. Direct terms  $(\langle \phi_i \phi_j | \delta(x x_k) | \phi_i \phi_j \rangle)$  survive and contribute  $|\phi_i(x)|^2$

2. Cross terms  $(\langle \phi_i \phi_j | \delta(x - x_k) | \phi_j \phi_i \rangle)$  vanish due to orthogonality 3. The normalization factor  $\frac{1}{N!}$  from  $\hat{A}$  exactly balances the number of equivalent permutations

Therefore, in general:

$$\rho(x) = \sum_{i} |\phi_i(x)|^2 \tag{13}$$

## **Double Occupation Case**

For doubly occupied spatial orbitals  $\varphi_j(r)$ :

$$\rho_{\alpha}(r) = \sum_{j} |\varphi_{j}(r)|^{2} = \rho_{\beta}(r)$$
(14)

The equality holds because each spatial orbital contributes identically to both  $\alpha$  and  $\beta$  spin densities.

## Problem 10 Part (a): Molecular Orbital Coefficients

#### **Given Information**

- MOs are linear combinations:  $\sigma_g = \chi_A C_{Ag} + \chi_B C_{Bg}$  and  $\sigma_u = \chi_A C_{Au} + \chi_B C_{Bu}$
- 1s functions are normalized:  $\langle \chi_A | \chi_A \rangle = \langle \chi_B | \chi_B \rangle = 1$
- Overlap integral:  $\langle \chi_A | \chi_B \rangle = S$

### Finding the MO Coefficients

Using symmetry properties:

- $\sigma_g$  is symmetric:  $C_{Ag} = C_{Bg}$
- $\sigma_u$  is antisymmetric:  $C_{Au} = -C_{Bu}$

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Let's impose normalization: For  $\sigma_g$ :

$$= \langle \sigma_g | \sigma_g \rangle$$
  
=  $(C_{Ag}^2 + C_{Bg}^2 + 2C_{Ag}C_{Bg}S)$   
=  $2C_{Ag}^2(1+S)$  (15)

Therefore:

$$C_{Ag} = C_{Bg} = \frac{1}{\sqrt{2(1+S)}}$$
(16)

For  $\sigma_u$ :

$$1 = \langle \sigma_u | \sigma_u \rangle$$
  
=  $(C_{Au}^2 + C_{Bu}^2 - 2C_{Au}C_{Bu}S)$   
=  $2C_{Au}^2(1-S)$  (17)

Therefore:

$$C_{Au} = -C_{Bu} = \frac{1}{\sqrt{2(1-S)}}$$
(18)

### Verifying Orthonormality

The coefficient matrix  ${\boldsymbol{\mathsf{C}}}$  is:

$$\mathbf{C} = \begin{pmatrix} \frac{1}{\sqrt{2(1+S)}} & \frac{1}{\sqrt{2(1-S)}} \\ \frac{1}{\sqrt{2(1+S)}} & -\frac{1}{\sqrt{2(1-S)}} \end{pmatrix}$$
(19)

Let's verify orthonormality:

$$\mathbf{C}^{\mathsf{T}} \begin{pmatrix} 1 & S \\ S & 1 \end{pmatrix} \mathbf{C} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$
(20)

## Problem 10 Part (b): Kinetic Energy Analysis

The kinetic energy expectation value is:

$$\langle \psi | \hat{T} | \psi \rangle = 2 \langle \sigma_g | \left( -\frac{1}{2} \nabla^2 \right) | \sigma_g \rangle + 2 \langle \sigma_u | \left( -\frac{1}{2} \nabla^2 \right) | \sigma_u \rangle$$
(21)

### **Qualitative Analysis of Kinetic Energy**





The  $\sigma_u$  orbital has a node and more spatial variation, leading to larger kinetic energy:

$$\langle \sigma_{u} | \left( -\frac{1}{2} \nabla^{2} \right) | \sigma_{u} \rangle > \langle \sigma_{g} | \left( -\frac{1}{2} \nabla^{2} \right) | \sigma_{g} \rangle \tag{22}$$

### Part (c): Coulomb and Exchange Energies

Using the properties of Slater determinant wavefunctions, we derive:

#### **Coulomb Energy**

$$J = 2 \langle \sigma_g \sigma_g | r_{12}^{-1} | \sigma_g \sigma_g \rangle + 4 \langle \sigma_g \sigma_u | r_{12}^{-1} | \sigma_g \sigma_u \rangle + 2 \langle \sigma_u \sigma_u | r_{12}^{-1} | \sigma_u \sigma_u \rangle$$
(23)

This represents:

- Direct interaction between electrons in  $\sigma_g$  orbitals
- Cross-interaction between  $\sigma_g$  and  $\sigma_u$  electrons
- Direct interaction between electrons in  $\sigma_u$  orbitals

#### **Exchange Energy**

The exchange terms arise from:

- Exchange between same-spin electrons in  $\sigma_g$  orbitals
- Exchange between same-spin electrons in  $\sigma_g$  and  $\sigma_u$  orbitals
- Exchange between same-spin electrons in  $\sigma_u$  orbitals

### **Physical Interpretation**

1. The molecular orbitals maintain proper symmetry ( $\sigma_g$  symmetric,  $\sigma_u$  antisymmetric)

2. The normalization factors depend on the atomic orbital overlap S, showing how bonding affects orbital structure

3. The higher kinetic energy of  $\sigma_u$  reflects the antibonding nature of this orbital

4. The Coulomb and exchange energies show both intra-orbital and interorbital electron interactions

5. The factor of 2 in various terms reflects the double occupation of orbitals